

Phase behavior of a family of continuous two-dimensional n -vector models with $n = 2, 3,$ and 4

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The presence of an unusual order disorder transition in the two dimensional XY model has been known since the pioneering works of Berezinskii¹ and Kosterlitz and Thouless². More generally, it has been found that similar transitions appear in lattice models of particular relevance in field theory and elementary-particle physics, such as the bidimensional RP^{n-1} (real projective space in n -dimensions).

In this work we investigate the phase behavior of a family of continuous bidimensional n -vector models (with $n = 2, 3,$ and 4) using Monte Carlo simulation, in which the explicit interaction is given by a hard sphere potential of diameter σ with embedded n -dimensional spins, i.e.

$$u(r_{ij}, \omega_i, \omega_j) = u_{HS}(r_{ij}) + u_{ang}(r_{ij}, \omega_i, \omega_j), \quad (1)$$

where $u_{HS}(r)$ is a hard sphere interaction, and the angular interaction is described by

$$u_{ang}(r, \omega_i, \omega_j) = -Ku_0(r) \frac{1}{n-1} [(\hat{s}_i \cdot \hat{s}_j)^2 n - 1] \quad (2)$$

with $\hat{s}_i = (s_{i_1}, s_{i_2}, \dots, s_{i_n})$ being a n -dimensional unit vector describing the orientation of the spin in particle i . In Eq. (2) the spin coupling is defined by

$$u_0(r) = \frac{e^{-\kappa(r-\sigma)}}{r/\sigma} - \frac{e^{-\kappa(R-\sigma)}}{R/\sigma}, \text{ for } \sigma < r < R, \quad (3)$$

$$= 0, \text{ for } r > R. \quad (4)$$

Here we have used both cluster algorithms in the line of Swendsen and Wang method³ with the aim of studying the order-disorder transition. First order phase changes (vapour-liquid equilibrium) has been studied using an algorithm based on Wang-Landau's method⁴. In all cases we detect the presence of a defect mediated order-disorder transition of the Berezinskii-Kosterlitz-Thouless type with critical temperatures that decrease with the spin dimensionality. Coupled with the order-disorder transition a gas-liquid equilibrium is found at low temperatures. Here one observes that the stability region of

the liquid phase shrinks with the growing spin dimensionality, in parallel with a decrease in magnitude of the angular averaged spin-spin interaction.

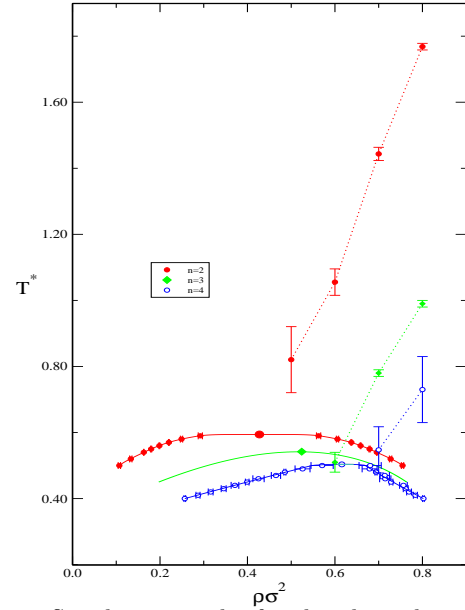


Figure 1. Simulation results for the phase diagram of the planar RP^{n-1} spin fluid for $n = 2, 3,$ and 4 . Symbols joined with dotted lines correspond to the states at which the BKT transition takes place and the line separates orientationally disordered states (lower densities) from states with quasi-long range order (higher density, quasi nematic states).

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